Digital resonance tuning of high- Q/V_m silicon photonic crystal nanocavities by atomic layer deposition

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We propose and demonstrate the digital resonance tuning of high- Q/V_m silicon photonic crystal nanocavities using a self-limiting atomic layer deposition technique. Control of resonances in discrete steps of 122 ± 18 pm/hafnium oxide atomic layer is achieved through this postfabrication process, nearly linear over a full 17 nm tuning range. The cavity Q is maintained in this perturbative process, and can reach up to its initial values of 49 000 or more. Our results are highly controllable, applicable to many material systems, and particularly critical to matching resonances and transitions involving mesoscopic optical cavities. © 2007 American Institute of Physics.

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Two-dimensional photonic crystal (2D PhC) slabs confine light by Bragg reflection in-plane and total internal reflection in the third dimension. Introduction of point and line defects into 2D PhC slabs create localized resonant cavities and PhC waveguides respectively, with full control of dispersion ab initio. Such defect cavities in high-index contrast materials, such as monolithic silicon, possess strong confinement with subwavelength modal volumes (V_m) at $\sim (\lambda/n)^3$, corresponding to high field intensities per photon for increased nonlinear and nonclassical interactions. Moreover, photonic crystal cavities with remarkable high quality factors (Q) (Refs. 1 and 2) have been achieved recently, now permitting nanosecond photon lifetimes for enhanced lightmatter interactions. The strong localization and long photon lifetimes in these high- Q/V_m photonic crystal nanocavities are strong candidates for enhanced nonlinear optical physics, such as optical bistability^{3–5} and Raman lasing,^{6,7} and cavity quantum electrodynamics.⁸

These applications require precise control of cavity resonances to achieve tuned spectral overlap between the cavity modes and the gain or emitter material for controlled lightmatter interactions. The cavity resonances is strongly dependent on the fabricated lattice constant a and the hole radius rof photonic crystals. Slight differences in the photonic crystal geometries will result in large differences (a few to tens of nanometers in wavelength) in the dispersion characteristics. Active tuning using thermal (with associated phonon broadening of quantum dots)⁹ or piezoelectric effects¹⁰ can be employed. A passive postfabrication tuning mechanism is particularly demanded, without external input power, to precisely align the designed resonant wavelengths. Specifically, wet chemical digital etching techniques¹¹ were recently developed for GaAs photonic crystal nanocavities, where the controlled blueshift of the cavity resonance was around 2-3 nm/cycle. Additionally, condensation of Xe (Ref. 12) or self-assembled monolayers (such as a 2 nm polypeptide monolayer)¹³ can be used, where a 3–5 nm cavity red shift/ monolayer was observed for the latter. To achieve hundreds of picometer cavity resonance tuning, thin films below 1 nm is needed to be etched or deposited.

Atomic layer deposition (ALD) is widely used for gate dielectric and capacitance memory applications due to its high dielectric constant, precise thickness control, and highly conformal properties. Several metal oxides such as aluminum oxide (Al₂O₃), hafnium oxide (HfO₂), and titanium dioxide (TiO₂) have been used in low temperature ALD thin film growth. All these materials have been widely used in optical coating applications with relatively high refractive indices [at 1.55 μ m, n=1.88 for HfO_2 , 14 n=1.57 for Al_2O_3 , 15 and n=2.18 for TiO_2 Ref. 16 and a wide band gap with low absorption from the near ultraviolet to the midinfrared. Recently, ALD has become a promising tool for the fabrication of high quality three-dimensional photonic crystals from inorganic (opal) and organic (patterned polymer) templates.^{17–20} Photonic band structure tuning in a 2D periodic lattice was also demonstrated, with 12% tuning range and 0.005% precision based on a deposition rate of 0.51 Å TiO₂/ALD cycle.²¹

In this letter, we report the postfabrication digital resonance tuning of high- Q/V_m silicon photonic crystal nanocavities using self-limiting atomic layer deposition of HfO₂.

The structure investigated is an air-bridged triangular lattice photonic crystal slab with silicon membrane thickness of 190 nm (t/a=0.4524) and air holes radii of 90 nm (r/a=0.2143), where the lattice period a=420 nm. High- Q/V_m nanocavities with five linearly aligned missing air holes (L5) are side coupled with photonic crystal waveguides, as shown in Fig. 1(a). The shift S_1 of two air holes at cavity edge is 0.02a, 0.06a, and 0.10a, respectively, for three different L5 nanocavities studied, in order to tune the radiation mode pattern for increasing the Q factors. The waveguide-to-cavity separation is five layers of holes. The devices were fabricated with the standard integrated circuit techniques in a silicon-on-inslator substrate. Next, optical lithography with AZ4620 photoresist was used to open a window in photonic crystal region, and 10 min HF BOE (6:1) was used to release the air-bridged structures. Samples were then cleaned using Piranha (H₂SO₄:H₂O₂, 3:1) solution for 5 min followed by HF BOE (6:1) solution dip for 30 s and de-ionized water rinse. This procedure results in \sim 6 Å of (O-H)-terminated silicon oxide on the surface of silicon airbridged photonic crystal slabs.²² All samples were exposed to UV generated ozone for 10 min to restore the hydrophilic

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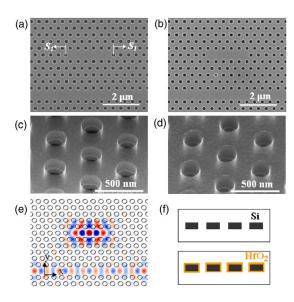


FIG. 1. (Color online) Top-view SEM images of airbridge L5 nanocavity with S_1 =0.02 α coupled with photonic crystal waveguides (a) before ALD and (b) after 140 ALD cycles of HfO₂. 45°-angle-view SEM images of air bridge photonic crystal slabs (c) before ALD and (d) after 140 ALD cycles of HfO₂. (e) 3D FDTD calculated electric field E_y profile of the high-Q mode supported in L5 nanocavities. (f) Deposition schematic cross section of the sample morphology.

character of surface immediately prior to HfO_2 deposition. Figure 1(a) shows the top-view scanning electron microscopy (SEM) image of air-bridged L5 cavity with S_1 =0.02a before ALD. Figure 1(e) shows the electric field E_y of the resonance mode midslab from three-dimensional (3D) finite-difference time-domain (FDTD) simulations, calculated using a freely available software package with subpixel smoothing for increased accuracy.²³

Thin films of amorphous HfO2 are deposited conformally on silicon air-bridged photonic crystal slabs by means of ALD at 150 °C. Films were deposited using tetrakis(diethylamido)hafnium (IV) [Hf(DEA)₄] and water (H₂O) vapor in alternating pulses with N₂ purge of the reaction chamber between pulses. Each deposition step includes 20 ALD cycles, and each cycle consists of Hf(DEA)₄ injection for 0.25 s, N_2 purge for 150 s, H_2O injection for 0.02 s, and N_2 purge for 200 s. The observed linear deposition rate is around 0.93 Å/cycle, which is about a monolayer of hafnium oxide. We note that lower substrate temperature down to 90 °C is possible with our machine at the expense of slow deposition rates. Figure 1(b) shows the top-view SEM image of L5 cavity after seven deposition steps, with the same magnification as in Fig. 1(a). Based on geometrical statistical analysis of high-resolution SEM images, the hole radius reduces from 92.84 ± 1.56 to 79.86 ± 2.66 nm. ²⁴ Figures 1(c) and 1(d) are the angled-view SEM images of air-bridged photonic crystal slabs before ALD and after seven deposition steps, respectively. The surface is still smooth enough to support high-Q modes for L5 nanocavities after HfO2 deposition. The thickness of photonic crystal slabs increases from 190 to 216 nm based on SEM estimates. These geometry changes agree well with the deposition schematic cross section of the sample morphology drawn in Fig. 1(f).

With slightly decreased r/a and increased t/a ratios in air-bridged photonic crystal slabs, the photonic band gap will shift to lower frequencies. In addition to a frequency shift, the photonic band gap also decreases from an 11.4% to a

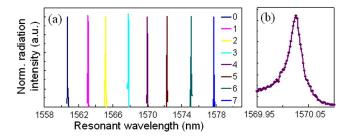


FIG. 2. (Color online) (a) Measured cavity resonances after each deposition step ("1" to "7" in legend; "0" is unperturbed) for L5 cavity with S_1 =0.02a. (b) Magnified resonance peak after the fourth deposition step, quality factor $Q \sim 49\,000$.

9.7% gap with a deposition of HfO_2 , computed using a freely available software package. This can be attributed to a lower-index contrast between the holes and the bulk dielectric. The resonant wavelength of L5 nanocavities will undergo a red shift.

For the measurement setup, a polarization controller and a lensed fiber are used to couple transverse-electric (TE) polarization light from tunable laser source into the waveguide. The vertical radiation from the top of nanocavities collected by a $40\times$ objective lens (numerical aperture of 0.65) and a 4× telescope was sent to the photodetector and lock-in amplifier to analyze the cavity resonances. In order to exclude optical nonlinear effects, low input power of 10 µW was coupled to the waveguide. Figure 2(a) plots the measured cavity resonances after each deposition step for L5 cavity with $S_1 = 0.02a$. Figure 2(b) magnifies the resonance peak after the fourth deposition step. The quality factor Q is estimated from the full width at half maximum and is \sim 49 000. From the 3D FDTD method, the Q factor and modal volume are calculated around 50 000 and ∼0.98 cubic wavelengths $[(\lambda/n)^3]$, respectively.

Figure 3(a) shows the tuned resonant wavelength scales linearly with the number of deposition steps for all three L5 cavities under investigation. Total resonant wavelength tuning range is around 17 nm with the current seven deposition steps. With more deposition steps, wider tuning range can be obtained. The 3D FDTD simulations [inset of Fig. 3(a)] show a linear shift in the resonant wavelength as expected from small perturbations, although there is more uncertainty in the simulations due to the high spatial resolution (approximately a few nanometers or less) required to capture this digital tuning. Figure 3(b) plots the resonant wavelength increment for each deposition step. An average wavelength red shift of 2.443 ± 0.359 nm is obtained for each step, which corresponds to a resonance shift of 122±18 pm/HfO₂ monolayer deposition. An oscillatory variation of the resonance shift is also observed, as shown in Fig. 3(b). This could be due to variations in the film deposition thickness, which is not perfectly the same in each step. In addition, we observe that the resonance increment itself increases slightly from 2.2 to 2.7 nm based on the linear curve fit. This is because, due to the conformality of ALD process, more dielectric material will be added relative to the previous step due to the expanded surface area, so that the resonance increment also slightly goes up, as illustrated in deposition schematics in Fig. 1(f).

With different deposition materials, the precision of resonant wavelength shift per ALD cycle can be changed. Single monolayer of HfO₂ induces an average 122 pm shift

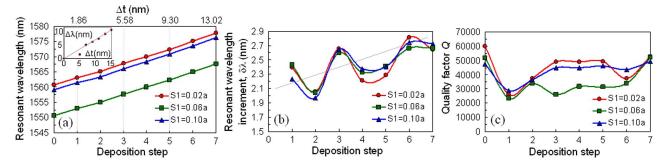


FIG. 3. (Color online) (a) The tuned resonant wavelength scales linearly with the number of deposition step for all three L5 cavities under investigation. Inset: 3D FDTD calculated wavelength shift $(\Delta\lambda)$ for increasing thicknesses (Δt) of HfO₂ deposited for all three cavities studied. (b) The wavelength increment $(\delta\lambda)$ for each deposition step. (c) The variation of quality factor Q with the number of deposition steps for all three L5 cavities.

(n=1.88 at 1.55 μ m, 0.93 Å/ALD cycle at 150 °C). From first-order perturbation estimate, a monolayer of TiO₂ [n=2.18 at 1.55 μ m, 0.5 Å/ALD cycle at 100 °C (Ref. 16)] can induce approximately 54 pm shift, while a monolayer of Al₂O₃ [n=1.57 at 1.55 μ m, 1 Å/ALD cycle at 100 °C (Ref. 15)] can generate approximately a 158 pm wavelength shift.

Figure 3(c) illustrates the variation of quality factor Q with the number of deposition steps for all three L5 cavities. After first deposition step, Q values drop almost by half for all cavities. This is because the ALD deposited film has a larger roughness initially, leading to more surface and air hole sidewall roughness scattering. With subsequent deposition steps, the conformal deposition gives a smoother film surface, permitting the Q values to recover back to almost their initial values. The Q values always maintain at least 20 000 or more during the deposition steps; this characteristic is also observed in our 3D FDTD simulations. This demonstrated shift in the resonance, while preserving the cavity Q, in response to a monolayer deposition also suggests these cavities as possible integrated sensors with pronounced responsivity to environmental conditions.

In summary, we have developed a technique for fine tuning the resonant wavelengths of high- Q/V_m silicon photonic crystal nanocavities digitally using ALD of HfO2 monolayers. The results demonstrate a nearly linear tuning across a range of around 17 nm. The tuning range is currently limited only by the number of deposition steps used in this study. The tuning precision is 122±18 pm/ALD cycle while preserving high quality factors of resonant modes in L5 photonic crystal nanocavities. With selective patterning, HfO₂ monolayers can be selectively deposited only within the nanocavity region using low-temperature ALD (Ref. 26) for even finer tuning control. The highly controlled, digital tuning of high-Q modes in silicon photonic crystal nanocavities allows for practical realization of optical devices involving multiple resonances and matching transitions between quantum dots and optical resonances for cavity quantum electrodynamics.

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- ¹B. S. Song, S. Noda, T. Asano, and Y. Akahane, Nat. Mater. **4**, 207 (2005).
- ²E. Kuramochi, M. Notomi, S. Mitsugi, A. Shinya, T. Tanabe, and T. Watanabe, Appl. Phys. Lett. **88**, 041112 (2006).
- ³P. E. Barclay, K. Srinivasan, and O. Painter, Opt. Express 13, 801 (2005).
 ⁴T. Tanabe, M. Notomi, S. Mitsugi, A. Shinya, and E. Kuramochi, Appl. Phys. Lett. 87, 151112 (2005).
- ⁵X. Yang, C. Husko, M. Yu, D.-L. Kwong, and C. W. Wong, Appl. Phys. Lett. **91**, 051113 (2007).
- ⁶H. Rong, S. Xu, Y. H. Kuo, V. Sih, O. Cohen, O. Raday, and M. Paniccia, Nat. Photonics 1, 232 (2007).
- ⁷X. Yang and C. W. Wong, Opt. Express **15**, 4763 (2007).
- ⁸R. Bose, X. Yang, R. Chatterjee, J. Gao, and C. W. Wong, Appl. Phys. Lett. **90**, 111117 (2007).
- ⁹E. Peter, P. Senellart, D. Martrou, A. Lemaître, J. Hours, J. M. Gérard, and J. Bloch, Phys. Rev. Lett. **95**, 067401 (2005).
- ¹⁰C. W. Wong, P. Rakich, S. G. Johnson, M. Qi, H. I. Smith, L. C. Kimerling, E. P. Ippen, Y.-B. Jeon, G. Barbastathis, and S.-G. Kim, Appl. Phys. Lett. 84, 1242 (2004).
- ¹¹K. Hennessy, A. Badolato, A. Tamboli, P. M. Petroff, E. Hu, M. Atatüre, J. Dreiser, and A. Imamoglu, Appl. Phys. Lett. 87, 021108 (2005).
- ¹²S. Mosor, J. Hendrickson, B. C. Richards, J. Sweet, G. Khitrova, H. M. Gibbs, T. Yoshie, A. Scherer, O. B. Shchekin, and D. G. Deppe, Appl. Phys. Lett. 87, 141105 (2005).
- ¹³S. Strauf, M. T. Rakher, I. Carmeli, K. Hennessy, C. Meier, A. Badolato, M. J. A. DeDood, P. M. Petroff, E. L. Hu, E. G. Gwinn, and D. Bouwmeester, Appl. Phys. Lett. 88, 043116 (2006).
- ¹⁴M. F. Al-Kuhaili, Opt. Mater. (Amsterdam, Neth.) **27**, 383 (2004).
- ¹⁵T. S. Eriksson, A. Hjortsberg, G. A. Niklasson, and C. G. Granqvist, Appl. Opt. **20**, 2742 (1981).
- ¹⁶K. Awazu, M. Fujimaki, X. Wang, A. Sai, and Y. Ohki, in *Nanoengineered Assemblies and Advanced Micro/Nanosystems*, MRS Symposia Proceedings No. 820 (Materials Research Society, Pittsburgh, 2004), p. R4.5.
- ¹⁷J. S. King, C. W. Neff, C. J. Summers, W. Park, S. Blomquist, E. Forsythe, and D. Morton, Appl. Phys. Lett. 83, 2566 (2003).
- ¹⁸A. Rugge, J. S. Becker, R. G. Gordon, and S. H. Tolbert, Nano Lett. 3, 1293 (2003).
- ¹⁹J. H. Lee, W. Leung, J. Ahn, T. Lee, I. S. Park, K. Constant, and K. M. Ho, Appl. Phys. Lett. **90**, 151101 (2007).
- ²⁰J. Y. Huang, X. D. Wang, and Z. L. Wang, Nano Lett. **6**, 2325 (2006).
- ²¹E. Graugnard, D. P. Gaillot, S. N. Dunham, C. W. Neff, T. Yamashita, and C. J. Summers, Appl. Phys. Lett. 89, 181108 (2006).
- ²²A. Deshpande, R. Inman, G. Jursich, and C. Takoudis, J. Vac. Sci. Technol. A 22, 2035 (2004).
- ²³A. Farjadpour, D. Roundy, A. Rodriguez, M. Ibanescu, P. Bermel, J. D. Joannopoulos, S. G. Johnson, and G. Burr, Opt. Lett. 31, 2972 (2006).
- ²⁴M. Skorobogatiy, G. Bégin, and A. Talneau, Opt. Express 13, 2487 (2005).
- ²⁵S. G. Johnson and J. D. Joannopoulos, Opt. Express **8**, 173 (2001).
- ²⁶M. J. Biercuk, D. J. Monsma, C. M. Marcus, J. S. Becker, and R. G. Gordon, Appl. Phys. Lett. 83, 2405 (2003).